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A Fast Computational Algorithm to Evaluate Large Transmit-arrays

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Abstract—Transmit-arrays (TAs) have become popular solutions for point-to-point communications due to their lowprofile and high-efficiency. However, simulation of these large and detailed structures demands a great amount of memory and time. Here, we propose a fast algorithm to estimate the performance of a beam-steering TA to accelerate its design optimization process. The results and the required resources of the algorithm are compared with the ones of Ansys HFSS. It is shown that this algorithm can fairly evaluate the key factors of the performance of a TA within couple of minutes with memory resource of only 1 GB.

Index Terms–Array antenna, trasnmit-arrays, reflectarrays, computational method.

I. INTRODUCTION

Passive transmit-arrays (TA) and reflect-arrays (RA) are very attractive solutions as gain enhancers in point-to-point and satellite communications [1]. These planar structures are usually populated with large number of unit-cells with dimensions smaller than the wavelength. The number of unitcells and the details of each unit-cell directly impact the amount of resources required to simulate such arrays. As a result, full wave simulations of high-gain planar arrays might take days even with high-performance computers.

Optimizing the design of these arrays and evaluating their performances before prototyping make their simulations an essential part of their design. However, these simulations can require a great amount of memory and time, which makes the optimization process time-consuming or even impossible based on the available resources. Therefore, a faster reliable technique rather than the full-wave simulation is necessary for this step. As a result, several methods have been proposed for RAs [8]-[12] and TAs [13]. Authors in [13] briefly describe a method to estimate only the main beam of a fixed-beam TA.

The literature yet lacks a general fast algorithm to evaluate both the main beam and the cross-polarization level. Here, we propose a general fast algorithm suited for TAs and/or RAs composed of any kind of unit-cells with even complicated shapes. These unit-cells are different in their geometries and in the incidence angles they receive the signal from the feed antenna. Based on the unit-cells' working principle, they can be categorized into two main groups: phase delay (PD) and phase rotation (PR). A set of PD cells includes a discrete number of cells with slightly different geometries, all with good transmission coefficient, in TA case, but slightly different phases shifts [2]-[5]. In the case of the PR cells, only one geometry for the PR cell is designed and then the cell is rotated by any angle from 0° to 180° to provide a continuous phase shift between 0° to 360° [6]-[9].

This case is important for various reasons including the fact that the scattered fields are at both sides of the TA and not just the primary feed side as in RAs. Secondly, the composing unit-cells have details much smaller than the wavelength that makes their simulation even more time and memory consuming.

II. DESCRIPTION OF THE ALGORITHM

To introduce the algorithm, we consider a TA similar to the one in [2] operating at 30GHz, with a rectangular aperture of size $D_v = 140$ mm × $D_x = 190$ mm, see Fig. 1. The size of the aperture is chosen to maximize the aperture efficiency for steering the beam from 15° to 50° zenith angles by in-plane shifting the feed from $x_{off} = +30mm$ to -27mm in Fig. 1. By in-plane rotation of the TA, 360° azimuth scanning can be obtained. The phase law of the TA is based on (1) where k_0 is the wavenumber at the design frequnecy and focal distance (F) is 100mm and the TA directs the beam to $\alpha_0=32.5^{\circ}$ when the feed is at the focal position of the TA. This α_0 is one of the design parameters of a beam-steering TA that can be optimized based on desired scanning range [2]. The TA is fed with standard 14.5dBi gain Ka-band rectangular horn with linear polarization. The performance of the TA is evaluated for both orthogonal linear polarizations by 90° rotation of the horn. Then the results are post-processed to be obtain the response of the TA to a synthesized circularly-polarized incident wave. Fig. 1 shows the configuration of the whole structure with the x-polarized horn antenna.

$$\varphi_{lens}(x,y) = k_0 \left(x \sin \alpha_0 - \sqrt{x^2 + y^2 + F^2} \right)$$
 (1)



Fig. 1 The configuration of the TA in front of the x-polarized horn antenna.

We populate the TA with the unit-cell similar to the one in [6]. The three metallic layers of the cell are separated by 0.508mm-thick RT Duroid 5880 (ε_r =2.2 and tan δ =0.0009). The exact dimensions of the unit-cell are presented in [14]. By rotating this unit-cell from 0° to 180°, it is possible to achieve up to 360° continuous phase shift in the transmitted wave, see Fig. 2. After populating the TA, it functions as shown schematically in Fig. 3. It collimates the orthogonal circular polarization of the feed antenna while it does not influence the cross-polarization component. Therefore, one component in the transmission is a high-gain plane-wave in the far-field while the other has a spherical phase front. The reverse stands true for reflections, see Fig. 3.



Fig. 2 The transmit-array unit-cell that provides phase shift from 0° to 360° by its rotation up to $180^{\circ}\!.$



Fig. 3 The behavior of the transmit-array and the scattered waves.

The amount of rotation angle of each cell and the incidence angle at which it receives the wave from the feed determines the cell scattering parameters (S-parameters). The incidence angle that different cells receive on the surface of the TA is depicted in Fig. 4. Therefore, the composing unitcells of this TA can be categorized based on their rotation angle and incidence angle at their surface.



Fig. 4 The incident angles from the feed positioned at the center of the TA and at focal distance of F=100mm across the surface of the TA in Fig. 1. Each pixel represents one unit-cell.

As the first step of the algorithm, we fit all the composing unit-cells into cases where the incidence angle changes (in this case) between 0° and maximum incidence angle on the TA, i.e. 45°, with 5° step ($N_{inc} = 10$) and the rotation angle changing between 0° and 170° with 10° step ($N_{cells} = 18$). The steps for these categories can be finer but the mentioned steps gave sufficient accuracy. This leads to $N_{inc} \times N_{cells} = 180$ different unit-cell simulations. The maximum amount of oblique incidence, here 45°, is obviously less for TAs with higher F/D.

Afterwards, we simulate the unit-cell for these 180 cases with periodic boundary condition and Floquet mode excitation in CST Microwave Studio [15], which can be done in any other full-wave solver. We then export all the Sparameters of the unit-cell for the 180 cases. This step is required only once for any type of unit-cell, and the obtained data base can be employed for different sizes of arrays and different positions of the feed antenna. Fig. 5 shows the magnitude and phase of the circular polarization Sparameters of the unit-cells populating the TA in Fig. 1 for different incident angles and different rotation angles.



Fig. 5 (a) Magnitude and (b) phase of the circular S-parameters of the unitcells rotated by different angles and excited by different incident angles on the TA in Fig. 1.

The second part of the proposed algorithm requires both the electric and magnetic fields of the primary feed at the focal distance of the TA. This can be exported from a fullwave simulation of the primary feed in any full-wave solver, calculated at the center of each cell. Afterwards, the tangential components of the scattered electric and magnetic waves can be obtained by multiplying the circular polarization S-parameters of the cells by the input electric and magnetic fields as in (2) and (3) respectively,

$$\begin{bmatrix} E_{S} \end{bmatrix} = \begin{bmatrix} e_{r}^{t} \\ e_{l}^{t} \\ e_{r}^{r} \\ e_{l}^{r} \end{bmatrix} = \begin{bmatrix} S_{21}^{rr} & S_{21}^{rl} \\ S_{21}^{lr} & S_{21}^{ll} \\ S_{11}^{rr} & S_{11}^{ll} \\ S_{11}^{lr} & S_{11}^{ll} \end{bmatrix} \begin{bmatrix} e_{r} \\ e_{l} \end{bmatrix}$$
(2)
$$\begin{bmatrix} H_{S} \end{bmatrix} = \begin{bmatrix} h_{r}^{t} \\ h_{l}^{t} \\ h_{r}^{r} \\ h_{l}^{r} \end{bmatrix} = \begin{bmatrix} S_{21}^{rr} & S_{21}^{rl} \\ S_{21}^{lr} & S_{21}^{ll} \\ S_{11}^{rr} & S_{11}^{rl} \\ S_{11}^{lr} & S_{11}^{rl} \\ S_{11}^{lr} & S_{11}^{rl} \end{bmatrix}$$
(3)

where e_r , e_l , h_r , and h_l are respectively the electric RHCP, electric LHCP, magnetic RHCP, and magnetic LHCP fields from the primary feed on the TA. Finally, the elements of E_s

and H_s represent the transmitted and reflected RHCP and LHCP waves from the TA that provide the near-fields to calculate the far-field patterns. Fig. 6 shows the magnitude and phase of the reflected and transmitted circular near-fields at the bottom and top surfaces of the TA, respectively.



Fig. 6 (a) Magnitude and (b) phase of the transmitted and reflected RHCP and LHCP electric near-fields at from the bottom and top surfaces of the TA for central position of the RHCP feed.

By calculating the electric and magnetic fields seperately based on (2) and (3), we avoid making a plane-wave approximation for the whole surface of the TA. Here, we only consider the plane wave approximation for each unit-cell. As can be seen from Fig. 6 (b), only the transmitted LHCP and the reflected RHCP fields have a planar phase front while the transmitted RHCP and the reflected LHCP fields have spherical phase fronts. It is noted that although we simulate the unit-cell with different rotation and incidence angles, we yet apply periodic boundary conditions in the unit-cell simulation that indicates the cell is surrounded by similar cells, which is not true in a beam-collimating TA.

As the last step of this method, the near-field data of the electric and magnetic fields are imported into our customized software [18] to calculate the far-field patterns. Any other

method that can calculate the far-field patterns from the near-field source can be employed at this step. Fig. 7 summarizes the steps of this algorithm.



Fig. 7 Algorithm to obtain far-field patterns of large transmit-arrays in Figure 3.

III. COMPARISON BETWEEN ALGORITHM AND FULL-WAVE SIMULATION

Here, we analyzed the introduced TA for three positions of the feed to steer the beam both by the proposed algorithm and Ansys HFSS [16]. It is worth mentioning that hybrid finite element boundary integral (FE-BI) method available by Ansys HFSS is employed in all full-wave simulations of this section in order to reduce the required time and memory [17]. The feed is placed at x_{off} = +30mm, 0mm, and -27mm, see Fig. 1. By mechanically shifting the feed to these positions, the beam based on (1) is directed at α_0 = 15°, 32°, and 50°, respectively.

In Fig. 8 (a), it is shown that both the full-wave simulation in HFSS and the proposed algorithm present similar beam direction and beam-width at each position of the feed. The maximum difference between the gain calculated by HFSS and by the new algorithm occurs for the beam directed at α_0 = 15°, which is 1.5dBi, while this difference reduces to 0.1dBi for the beam directed at α_0 = 50°. Moreover, the scan loss over this scan range is 4.1dBi calculated by the algorithm which is only 1.4dBi more than the one obtained by HFSS.

Fig. 8 (b) presents the far-field patterns of the crosspolarization component using both tools. As expected and explained earlier, this TA has no phase correcting effect on the cross polarization component, and consequently the cross polarization far-field is not pointed at the main beam direction and has a maximum in vicinity of $\theta=0^{\circ}$. Fig. 8 (b) shows that the proposed algorithm correctly estimates the overall shape of the cross component patterns although the cross polarization gain calculated by it is 3.4dBi higher than the one obtained with HFSS.



Fig. 8 Far-field patterns of the TA using full-wave simulation in Ansys HFSS [16] and using the algorithm.

The comparison with full wave Ansys HFSS calculations proves that the proposed algorithm can correctly estimate the overall performance of a large TA. Therefore, one can easily use this tool at the beginning of a design process to determine the general parameters of a TA such as F/D, D_x and D_y in case of a rectangular aperture, and α_0 . It allows fast estimation of main beam direction, beam-width, gain of the TA as well as cross polarization level and scanning loss. Furthermore, this algorithm provides a way for a designer to gain confidence in her/his design before simulating the whole structure using an accurate but time-consuming full-wave simulation software. As a result, the full-wave simulation can be done only once as the last step before moving to prototyping. Table I compares for the discussed TA the resources such as time and memory required for the simulation of each polarization and position of the feed using HFSS and the proposed algorithm.

TABLE I COMPARISON BETWEEN REQUIRED RESOURCES FOR HFSS AND THE ALGORITHM

Tool	Time	Memory
Ansys HFSS	17h	108 GB
Proposed Algorithm	2m 10s	1 GB

IV. CONCLUSION

A fast algorithm is presented to estimate the key performance characteristics of a beam-steering transmit-array such as the main beam direction, beam-width, gain, cross polarization level, and scan loss. It was shown that this algorithm can be employed even for non-collimating TAs where the scattered field does not have a planar phase front. The accuracy and the required resources of the algorithm are compared to a full-wave simulation software, i.e. HFSS. It is proved that, after a one-time full-wave analysis of a finite set of individual cells with different incidence angles, the proposed algorithm can expedite the optimization of large planar arrays before prototyping. A 500× factor was obtained for the presented example. Last but not the least, this algorithm can also be used for other types of transmit-array cells such as PD cells in [2]-[5] as well as reflect-arrays cells.

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